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### Key Points:

- A multiple-flyby mission to Europa can recover key geophysical parameters
- Laser altimetry can uniquely and accurately recover the global shape of Europa
- Laser altimetry enables the recovery of  $h_2$  to constrain the ice shell thickness

### Supporting Information:

- Figure S1

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## Simulated recovery of Europa's global shape and tidal Love numbers from altimetry and radio tracking during a dedicated flyby tour

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**Abstract** The fundamental scientific objectives for future spacecraft exploration of Jupiter's moon Europa include confirmation of the existence of subsurface ocean beneath the surface ice shell and constraints on the physical properties of the ocean. Here we conduct a comprehensive simulation of a multiple-flyby mission. We demonstrate that radio tracking data can provide an estimate of the gravitational tidal Love number  $k_2$  with sufficient precision to confirm the presence of a liquid layer. We further show that a capable long-range laser altimeter can improve determination of the spacecraft position, improve the  $k_2$  determination ( $<1\%$  error), and enable the estimation of the planetary shape and Love number  $h_2$  (3–4% error), which is directly related to the amplitude of the surface tidal deformation. These measurements, in addition to the global shape accurately constrained by the long altimetric profiles, can yield further constraints on the interior structure of Europa.

## 1. Background and Motivation

Despite the high-radiation environment of Jupiter's radiation belts, its moon Europa is one of the Solar System bodies with greatest potential for life, due to the presence of a subsurface ocean, nutrient exchange with the silicate mantle on its floor, and an energy source in the form of tidal heating. Dedicated exploration remains a high priority for the scientific community [National Research Council, 2011]. The thickness of the moon's ice shell is critical with regard to understanding the possible exchange of material between the ocean and surface, and the scientific potential of future lander exploration. Both the long-wavelength shape of Europa and its tidal response to Jupiter's time-variable gravitational forcing can help address this geophysical objective.

After the release of the Planetary Science Decadal Survey, various mission architectures were studied, including a multiple-flyby scenario. Because this concept did not include a laser altimeter, key geophysical parameters would not be measured, in particular the global shape and the surface tidal deformation. The lack of these measurements would make the characterization of the ice shell thickness difficult, because of the large trade-off between thickness and viscosity of the ice [Wahr *et al.*, 2006].

Analogous to the Galileo and Cassini missions, the spacecraft would conduct a series of flybys over the course of several years while in orbit around Jupiter, but with an almost exclusive focus on Europa. We used the notional trajectory 13F7-A21 developed by the Europa Pre-Project, optimized for spatial and phase coverage and surface illumination. It consists of 45 flybys between March 2029 and October 2031. We assumed that medium-gain antennas allow continuous tracking from Earth via the Deep Space Network when in proximity of Europa, and that the spacecraft is tracked in X band ( $\sim 8$  GHz) over the 4 h surrounding the closest approach, in addition to the typical 8 h long daily tracking passes. The plasma noise is the main contributor to the radio error budget, and we compute the observation noise level at 60 s sampling from the Sun-Probe-Earth angle of each flyby, following *less et al.* [2014]. The radio data are thus significantly degraded for some flybys (e.g., #39 in June 2031).

*Park et al.* [2011] used an earlier version of the Europa flyby tour trajectory to perform a simulation through covariance analysis to assess the recoverability of Europa's tidal Love number  $k_2$ . They found that a higher frequency Ka band ( $\sim 32$  GHz) system could yield an uncertainty of  $\sim 10\%$  (0.05), sufficient to detect the presence of a subsurface ocean. In this work, we performed a comprehensive simulation

of the complete Europa Clipper flyby tour with both radiometric and altimetric measurements, in an attempt to retrieve more key geophysical parameters, in particular the surface displacements not detectable by gravity alone. With six flybys by Cassini, *Jess et al.* [2012] recovered the tidal Love number for Titan to  $\sim 12\%$  accuracy.

This paper is organized as follows. Section 2 details the assumptions about the laser altimeter, the altimetric data quality, and sampling used in the simulation. Section 3 describes the simulation setup, including our methods, assumptions for a priori models, and initial errors, as well as the data coverage and noise sources. In sections 4 and 5, we present the results of a covariance analysis and the full simulation, respectively. Our results quantitatively demonstrate the quality of the recovery of the long-wavelength shape from altimetry and that of the low-degree gravity and Love number parameters from combined radio tracking and altimetric crossovers.

## 2. Laster Altimeter

### 2.1. Instrument Description

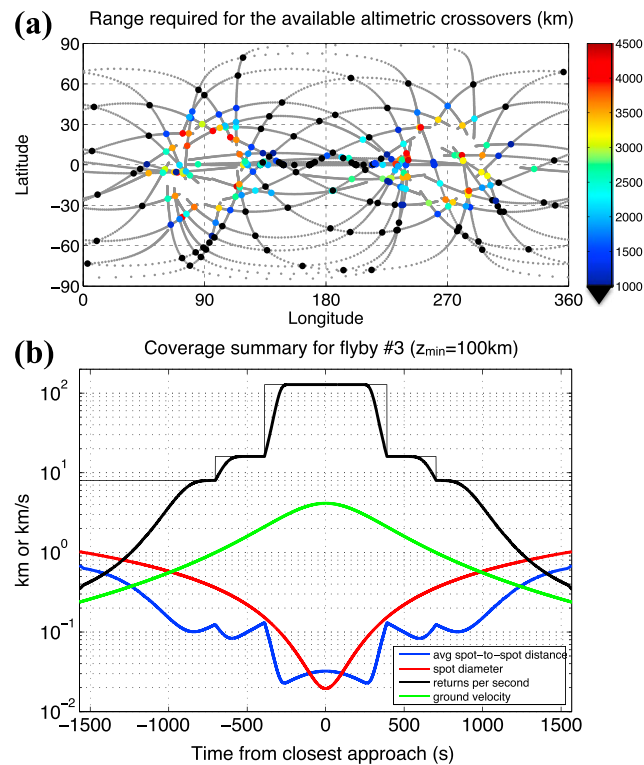
Laser altimeters measure topography with high accuracy from the time of flight of a short laser pulse reflected on the surface. These ranges are also strong geodetic constraints when used in Precise Orbit Determination (POD). Like other types of active instruments, laser altimeters have a maximum operating range (typically  $<400$  km) and thus have generally been limited to orbital missions [Smith et al., 1997, 2001; Zuber et al., 1997; Araki et al., 2009; Smith et al., 2010]. The Mercury Laser Altimeter (MLA) detected ground returns up to  $\sim 1800$  km altitude [Zuber et al., 2012] in a highly eccentric orbit and was the basis for the instrument concept for the simulation.

The simulated altimeter is a high-performance and flexible instrument that can operate at variable pulse energy and firing frequency to best match the changing observation geometry over the course of a flyby. Operations would start at  $>4500$  km with 8 Hz pulses, which dramatically improves the spatial coverage in terms of both quantity and uniformity (Figure 1a). A typical flyby yields  $\sim 100,000$  altimetric returns over a  $>150^\circ$  long arc. The precise determination of the gravity field by radio tracking and of the shape of Europa by an altimeter can place constraints on the ice shell compensation and thickness variations [Nimmo et al., 2007, 2011; Mitri et al., 2014] but also support the calibration of camera images and stereographic reconstructions [Preusker et al., 2011; Scholten et al., 2012].

The operating frequency is increased as needed to obtain the highest number of ground returns, from 8 Hz at  $\sim 4500$  km range up to a maximum frequency of 128 Hz below  $\sim 1000$  km altitude. As shown in Figure 1b, the spot size increases and the ground speed decreases with altitude, resulting in nearly contiguous spots for the majority of the flyby duration, thanks to the variable operating laser firing frequency. Even at closest approach, typically 100 km, the spots ( $\sim 10$  m) are separated by only small gaps ( $\sim 20$  m). Except in specific locations within chaos-like regions ( $<20\%$  surface area) [Figueredo and Greeley, 2003] where extreme slopes can occur, this sampling is sufficient to prevent significant interpolation errors compared to the tidal signal amplitude ( $\sim 30$  m) [Sotin et al., 2009]. Indeed, with a maximum typical root-mean-square (RMS) surface slope of  $10^\circ$ , the worst-case interpolation error is  $\sim 80$  cm, which is of the same order as the measurement noise ( $<1$  m including detector response and timing error, similar to MLA at long range) [Zuber et al., 2012]. With large spot sizes, the range can be somewhat biased by the spot-scale slopes; this effect is on the order of a few meters at most, but is partially compensated by the smoother topography sampled over larger length scales, due to the fractal nature of planetary surface topography. The reconstructed pointing knowledge of the spacecraft is assumed to be the same as the beam divergence of the laser,  $100 \mu\text{rad}$ .

### 2.2. Shape Recovery

The geodetic accuracy of the laser altimetric ranges yields the global shape of Europa. The long-range capability is especially important as it provides near-global coverage with few regional gaps (Figure 1a). In order to illustrate the quality of the recovery of the low-degree shape, we performed constrained inversions to estimate the spherical harmonics coefficients that describe the global shape to degree 32 (spatial block size  $\sim 150$  km), given a topographic model and the ground tracks of the notional Europa Clipper mission. We used several a priori shapes (Figure S1).



**Figure 1.** (a) Spacecraft ground tracks (up to an altitude of 4500 km; sampled every 10 s) and the resulting crossover locations. Color indicates the altitude of the higher of the two intersecting tracks (in km) and highlights that the  $90^\circ/270^\circ$  regions are not well sampled with a  $<1000$  km range capability. (b) Key parameters during a typical flyby ( $\sim 100$  km closest approach). The commanded firing frequency of the instrument (thin) and the effective measurement rate (thick) are in black. As the spacecraft approaches Europa, its ground speed (green) increases and the laser footprint shrinks (red). Given the measurement rate, the average spot-to-spot distance (blue) can be computed. During the majority of the flyby, the spots are overlapping or near contiguous (i.e., the red curve is above the blue one).

spacecraft missions [e.g., Lemoine *et al.*, 1997, 2001, 2013; Mazarico *et al.*, 2014]. Detailed force and measurement models are used to iteratively integrate the spacecraft trajectory and minimize the difference between the computed and measured observables [cf. Lemoine *et al.*, 2013], a process called Precise Orbit Determination (POD). Lemoine *et al.* [2001] describe the altimetric crossover measurement type and its use with Mars Orbiter Laser Altimeter [Zuber *et al.*, 1992] data for POD of the Mars Global Surveyor spacecraft. For each of the 45 flybys, we focused on 24 h centered on the closest approach. The integration of the spacecraft trajectory is initiated  $\sim 200,000$  km away from Europa, outside of its sphere of influence.

To construct the truth gravity field, we combined the degree 2 gravity spherical harmonics of Europa [Anderson *et al.*, 1998], the degrees 3 and 4 from estimates of Titan's [Joss *et al.*, 2010], and the higher degrees (to 50) from a Mars field [Konopliv *et al.*, 2011] scaled appropriately to simulate the silicate mantle contribution [Pauer *et al.*, 2010]. The tidal parameters are  $k_2 = 0.257$  and  $h_2 = 1.20$ , which translate to  $\sim 30$  m radial tidal deformation and  $\sim 35$  km ice thickness [Wahr *et al.*, 2006].

A 20 km wide 20 m/pixel stereo-derived digital elevation model (DEM) of Europa in a ridged plains region imaged by Galileo during the E4 flyby [Head *et al.*, 1999; Nimmo and Schenk, 2008] was used as representative of the global surface topography (Figure 2a). We added short-wavelength fractal noise to obtain a topography model of higher resolution, commensurate with the smallest altimeter footprints

Table 1 shows that the coverage obtained by the altimeter yields substantial improvements over what a shorter-range instrument operating between 400 and 1000 km could achieve (Figure S1). The  $C_{20}$ ,  $C_{22}$ , and  $S_{22}$  shape coefficients define the flattening, equatorial ellipticity, and orientation, improve by a factor of 5–10 due to the longer ground tracks. In combination with the low-degree gravity coefficients, the determination of the low-degree shape to high accuracy enables several types of studies, such as its orientation with respect to the (gravitational) figure, the degree of compensation in the ice shell [Ojakangas and Stevenson, 1989; Hemingway *et al.*, 2013; Mitri *et al.*, 2014; Lefèvre *et al.*, 2014], and indirect constraint on the heat flow. For example, the degree 4 terms, also well determined with the greater altimeter coverage, are important because they are not expected to arise from tidal or rotational effects but could be large depending on the heat loss mechanism of the ice shell (conductive versus convective) [Nimmo *et al.*, 2007].

### 3. Simulation Setup

We use the NASA Goddard Space Flight Center's GEODYN II orbit determination and geodetic parameter estimation software, which has been used for geophysical studies on numerous

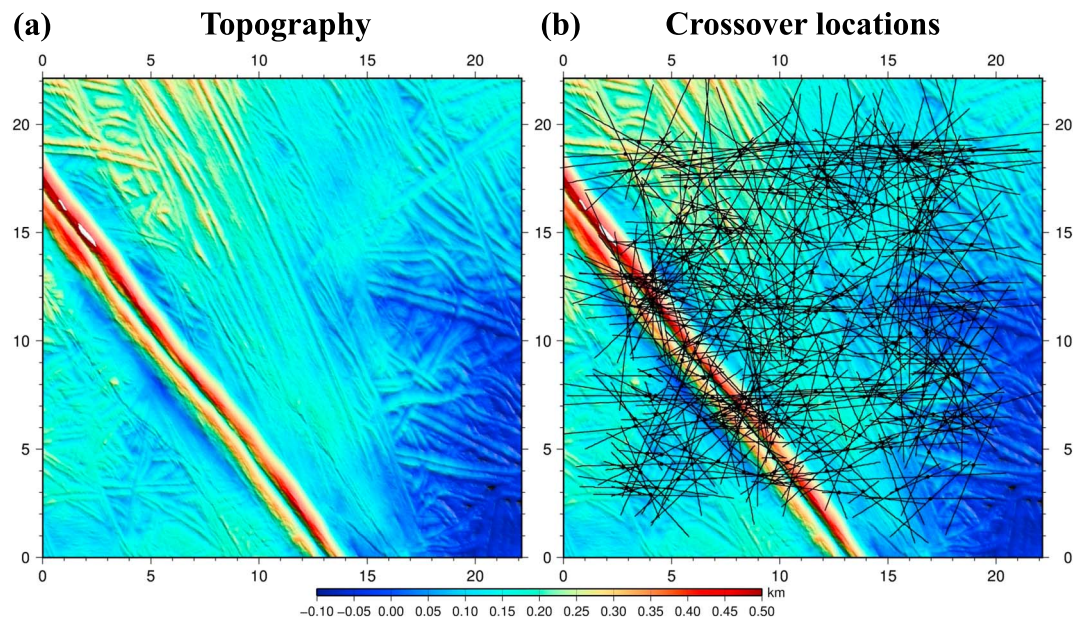
**Table 1.** Values and Errors (in m) of the  $C_{20}$ ,  $C_{22}$ ,  $S_{22}$ , and  $C_{40}$  Spherical Harmonics Coefficients Recovered From Various A Priori Shapes With Different Ground Track Coverages<sup>a</sup>

Shape	$C_{lm}$ or $S_{lm}$	Truth		$z < 4500$ km		$400 \text{ km} < z < 1000$ km		
		Value		Value		Error		Error
		(m)	(m)	(m)	(%)	(m)	(m)	
Moon	$C_{20}$	−668.55	−651.63	16.92	2.5	−524.16	144.39	21.6
	$C_{22}$	109.63	141.25	31.62	28.8	287.3	177.67	162.1
	$S_{22}$	383.11	374.8	8.31	2.2	360.19	22.92	6.0
	$C_{40}$	215.89	239.43	23.54	10.9	70.15	145.74	67.5
Hydrostatic shape based on rescaled Mimas	$C_{20}$	−278.48	−278.31	0.17	0.1	−257.94	20.54	7.4
	$C_{22}$	−83.69	−82.61	1.08	1.3	−67.46	16.23	19.4
	$S_{22}$	−0.74	3.55	4.29	579.7	1.57	2.31	312.2
	$C_{40}$	13.87	16.94	3.07	22.1	10.7	3.17	22.9
<i>Nimmo et al.</i> [2007]	$C_{20}$	−45	−44.51	0.49	1.1	−26.73	18.27	40.6
	$C_{22}$	−103.32	−102.77	0.55	0.5	−86.79	16.53	16.0
	$S_{22}$	0	0.17	0.17	−	0.86	0.86	−
	$C_{40}$	21	20.63	0.37	1.8	12.5	8.50	40.5

<sup>a</sup>Maps of the a priori and recovered degree 10 expansions are shown in Figure S1. Numbers in italics are large due to the small a priori values.

(10 m/pixel). The RMS of these short-wavelength and midwavelength perturbations is 5 and 50 m, respectively, which produces a realistic slope distribution [Schenk, 2009].

To ensure the robustness of the simulation results, we strongly perturbed the truth parameters. Because we found that degree 4 coefficients are not well determined with the flyby geometry (signal-to-noise ratio  $\sim 1$ ), we only adjusted the coefficients of degrees 2 and 3, which introduces realistic omission error. The imposed initial error on the  $C_{20}$  and  $C_{22}$  coefficients is 10 times the uncertainty given by Anderson *et al.* [1998], while the other coefficients are initially zero. The initial  $k_2$  and  $h_2$  are 0.107 ( $\sim 60\%$  error) and 0.60 (50% error), respectively. The scale factor for the solar radiation (modeled with a plate model of the spacecraft and nadir orientation) is perturbed by 20% but held fixed because of high correlation with the state. Future simulation efforts will include the perturbing effects of the Radioisotope Thermoelectric Generators or of

**Figure 2.** (a) Typical Europa topography used to compute the artificial altimetric ranges around the times of crossovers. (b) To capture some of the variability in terrain around Europa, each crossover center was randomly placed in the DEM, as indicated by the black dots and intersecting short segments. Axes are in kilometers.



**Table 2.** Values and Errors of the Tidal Love Numbers and Gravity  $C_{20}$ ,  $C_{22}$ ,  $S_{22}$ , and  $C_{30}$  Spherical Harmonics Coefficients Recovered From the Various Simulation Cases

	Truth	Covariance Analysis			Radio Only			Radio + Altimetry		
	Value	Value	Error	(%)	Value	Error	(%)	Value	Error	(%)
$k_2$	0.257	0.2572	0.0002	0.08	0.2495	0.0075	2.93	0.2546	0.0024	0.94
$h_2$	1.2	1.182	0.018	1.50	-	-	-	1.1660	0.0340	2.83
$C_{20}$	-4.3550E-04	-4.3547E-04	3.0000E-08	0.01	-4.3707E-04	1.5675E-06	0.36	-4.3672E-04	1.2165E-06	0.28
$C_{22}$	1.3100E-04	1.3097E-04	3.0000E-08	0.02	1.3228E-04	1.2831E-06	0.98	1.3158E-04	5.7729E-07	0.44
$S_{22}$	-1.1900E-05	-1.1896E-05	4.0000E-09	0.03	-1.1998E-05	9.7800E-08	0.82	-1.2193E-05	2.9255E-07	2.46
$C_{30}$	2.5700E-05	2.5723E-05	2.3000E-08	0.09	2.5889E-05	1.8917E-07	0.74	2.5887E-05	1.8713E-07	0.73

the larger area of the solar panels if such a spacecraft configuration is selected. The spacecraft position and velocity before each flyby are initially perturbed in each direction by 1 km and 1 mm/s, respectively. Planetary orientation and ephemeris will be significantly improved as the mission progresses, and we conducted alternate simulations to verify our sensitivity to initial errors in these parameters.

To create the simulated tracking and altimetric data for each flyby, we used GEODYN II to obtain the state of the spacecraft which, when propagated with our modeling assumptions, resulted in the best match to the Europa Clipper trajectory. Along these hyperbolic paths, Doppler observables were computed at 10 s intervals, and we added Gaussian noise with an RMS computed from the Sun-Probe-Earth (SPE) angle (typically 0.06–0.11 mm/s, with a few flybys with noise greater than 1 mm/s). Altimetric ranges to the surface were computed with nadir orientation and all necessary relativistic corrections. The average measurement frequency was derived from spacecraft altitude and the altimeter performance (Figure 1b). For each crossover track, we smoothed the surface shape model in order to match its effective resolution to the laser footprint size, and we placed the crossover point randomly in the DEM grid so that the crossovers sample a variety of terrain (Figure 2b). We then added ~1 m noise to the altimetric ranges and grouped each pair of short segments around the precomputed crossover locations and times, for use in GEODYN II as a crossover measurement. The altimeter spot location was assumed to be within the beam divergence of the laser and the reconstructed knowledge of the spacecraft attitude (100  $\mu$ rad).

#### 4. Covariance Analysis

For the radio science simulation, we first performed a covariance analysis, similar to *Park et al.* [2011], but note here that the trajectory we used (13F7-A21) had been optimized, in part to provide better tidal phase sampling. Partial derivatives are computed along the truth trajectories and with no perturbation of the model parameters. The RMS of the Doppler observations is of course at noise level, and the RMS of the crossover discrepancies is ~8 m. We discarded the crossovers with large discrepancies (>100 m), which can occur due to the sampling and the local terrain (<20% of crossovers).

The least squares solution provides both adjustments (i.e., errors) and formal uncertainties. We find that the formal uncertainties from the covariance analysis are 3 to 20 times smaller than the errors. Part of the reason the formal uncertainties in the spacecraft states might be underestimated is because we do not adjust the solar radiation scale factors. Still, the covariance results (Table 2) indicate a recovery of  $k_2$  to 0.08% (0.00022) and of  $h_2$  to 1.5% (0.018). The low-degree gravity coefficients show errors of ~0.01% for degree 2 and ~0.1% at degree 3. The covariance analysis does not properly account for the concealed correlations and trades between parameters that do not exist around the global minimum.

#### 5. Full Simulation

##### 5.1. Radio Tracking Only

More realistic error estimates can be obtained with a full simulation, where POD is performed on noisy data starting from imperfect models. This is especially important in a flyby tour, where the geometry is not as robust as an orbital configuration because the dynamics are weaker (hyperbolic trajectory versus multiple short-period orbits) and long gaps exist between the flybys.

While the altimetric crossovers are a powerful measurement of the surface deformation, we found that POD should first be performed with the radiometric tracking data alone. Indeed, with large perturbations, the initial geolocation errors are such that crossovers can become degenerate due to insufficient data selection (we typically selected 50 points centered on the expected intersection, which is only a short amount of time at 128 Hz) or due to crossover discrepancies so large that they destabilize the iterative least squares process because of the inherent nonlinearity of the altimetric profiles.

To remediate this issue and realistically follow the mission scenario, we processed the radio tracking data with a batch-sequential approach, similar to *Genova et al.* [2013] with data from the MErcury Surface, Space ENvironment, GEochemistry, and Ranging (MESSENGER) mission. The first flyby is processed alone; updated arc and global parameters are used as a new a priori to process the first and second flybys, and so on. Once all 45 flybys are processed, we obtained the radio-only solution (Table 2). A few of the flybys do not contribute much to the solution, because of high radio noise due to low SPE angles. In particular, the error in the initial state of flyby #39 is very large ( $\sim 7$  km).

The formal uncertainties obtained were once again optimistic, by a factor of 5–20 compared to the actual errors. The errors on the degree 2 and 3 gravity coefficients (Table 2) are typically  $\sim 1\%$  and  $\sim 10\%$ , respectively, except for the zonal terms which are better determined. The tidal Love number  $k_2$  is obtained to  $\sim 3\%$  accuracy (0.0075), which would enable the definitive confirmation of the presence of a subsurface ocean. If not all the flybys are available for radio tracking, the recovery accuracy will degrade accordingly. Simulations carried out with 10 and 20 flybys indicate that the  $k_2$  parameter can be estimated to  $\sim 16\%$  and  $\sim 7\%$  accuracy, respectively, consistent with the results of *less et al.* [2012] at Titan with Cassini. With 35 flybys, the  $k_2$  can be recovered to similar accuracy as with the complete flyby tour.

## 5.2. Radio Tracking and Altimetry

To account for instrument noise and interpolation error, we weighted each crossover measurement at a nominal 10 m. Although there are fewer than 300 crossovers, this value gives significant weight to the crossovers relative to the radio data. Thus, we used the full covariance matrix of the radio-only solution to provide an a priori constraint on the initial states of each flyby when processing altimetry, to prevent spurious noise (such as topographic interpolation errors) from destabilizing the solution. We scaled each  $6 \times 6$  state covariance matrix by a factor of 50 to be commensurate with the actual error level. In general, we find that the Z direction (in J2000) is the least well determined, so the states in the other two directions are not allowed to change significantly. The arc parameters and global tidal and  $L = 2$  and 3 gravity coefficients are iteratively adjusted with the complete data set. The final RMS of the crossovers is  $\sim 10.3$  m, higher than the covariance analysis but adequate given the tidal signal amplitude.

We found that, through the use of crossovers, the superior observation geometry of the great majority of the flybys helps better determine the states of the weaker flybys (from a radio perspective). For instance, the error in the state of flyby #39 is reduced by more than 1 order of magnitude to 500 m. The benefits of dynamics with the POD of the spacecraft for entire orbits around Jupiter would likely provide further improvement, but this is outside of the scope of this work.

The addition of altimetry also improves the gravitational parameters, in particular  $k_2$ . The error is reduced by a factor of 3, to 1% (0.0026), a level of accuracy amenable to use in interior models of Europa as has been performed on the Moon by *Williams et al.* [2014] with data from the Gravity Recovery and Interior Laboratory (GRAIL) mission.

The  $h_2$  Love number, not measurable from radio data alone, is determined to better than 3% (0.034), which would provide some constraints on the ice shell thickness when combined with  $k_2$ . *Wahr et al.* [2006] showed that the trade-off between ice shell thickness and viscosity is largely eliminated when considering the parameter  $1 + k_2 - h_2$ . In our case, the error in  $h_2$  dominates and results in a  $\sim 20$  km uncertainty in ice shell thickness. Future work, such as a joint inversion of  $h_2$  and  $k_2$ , may help reduce it. However, such a determination could provide a useful constraint in the absence of any other observation of the depth of the crust-ocean interface. A reduction in the number of flybys tracked by radio near close approach can degrade the recovery of  $h_2$ . With only 35 and 20 flybys tracked, the uncertainty in  $h_2$  increases to  $\sim 7\%$  and  $\sim 36\%$ , respectively.

In order to assess the sensitivity of our results to surface roughness, we varied the amplitude of the fractal noise added to the topographic maps, either short wavelength or midwavelength, by  $\pm 50\%$ . In some of

these cases, which are not expected from our current knowledge of Europa's topographic range [Nimmo *et al.*, 2007], we found that the error in  $h_2$  can increase, up to ~6% accuracy (0.075). We also used a DEM of Arbela Sulcus on Ganymede [Giese *et al.*, 2001] instead of the Europa relief map, and we obtained very similar results (1% and 3.6% for  $k_2$  and  $h_2$ , respectively).

We conducted another set of simulations for the case of an altimeter of limited range (1000 km). The measurement noise was increased to 10 m to reflect the larger uncertainty in geolocation. More significantly, we limited the crossovers to only those occurring when the spacecraft altitude is between 400 and 1000 km. The resulting errors in  $h_2$  were typically of the order to 25% (~0.3), too large to significantly contribute to detailed geophysical studies.

Although we anticipate the quality of the orientation parameters of Europa to be improved by the time a mission reaches it, we performed a simulation where the right ascension (RA) and declination (DEC) of the pole position were both perturbed by  $0.5^\circ$  (several times the current estimated quality) [Margot *et al.*, 2013]. After a sequential radio-only estimation and the addition of altimetric data, we found that the RA and DEC could be recovered to  $\sim 0.003$  and  $\sim 0.02^\circ$ , respectively. This is of sufficient quality to verify theoretical predictions of obliquity related to the presence of a global liquid layer, such as Bills *et al.* [2009], Baland *et al.* [2012], and Van Hoolst *et al.* [2013].

Similarly, we perturbed the Europa ephemeris by  $\sim 250$  m in each direction through Set III parameters [Brouwer and Clemence, 1961]. These perturbations are recovered to better than 50 m, and like Park *et al.* [2011], we find that the geophysical parameters are not affected by the recovery.

## 6. Summary

We showed that altimetric data from a capable instrument can substantially further the geophysical objectives of the Europa Clipper mission. The global coverage enabled by its long-range capability would deliver the global shape of Europa with high accuracy, which would constrain the ice shell thickness variations and thermal loss mechanism. We conducted a comprehensive simulation with radiometric tracking and altimetric data over the 2.5 year Europa flyby tour. Compared to the radio-only case, this combination improves the estimate of the tidal Love number  $k_2$  by a factor of 3, to  $<1\%$  accuracy, makes possible a measurement of  $h_2$ , to 3–4% accuracy, and improves the overall knowledge of the spacecraft position. This demonstrates a laser altimeter on board a Europa flyby mission will strengthen and enhance the geophysical investigation.

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